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Steinhardt, Jeong, Saitoh, Tanaka, Abe & Tsai [Nature (London) (1998), 396, 55–57] have demonstrated that the structure of decagonal Al–Ni–Co can be built from overlapping clusters of a single type. The structure arises from a decoration of the decagons of a Gummelt covering. The unit (essentially a decagonal prism) was called by Steinhardt et al. a ‘quasi unit cell’. In this work, a classification scheme is proposed for ‘G patterns’ – quasiperiodic patterns obtained by decorating a decagonal quasi unit cell. The classification makes use of the fact that G patterns can also be derived from decoration of a tiling. The tiles are analogues, for decagonal quasiperiodic patterns, of the ‘asymmetric units’ of a periodic pattern; they provide a simple mode of description and classification of the ‘Gummelt-type structures’. Four existing models for decagonal phases are considered from this viewpoint.

1. Introduction

A decagonal quasicrystalline phase can be thought of as a stack of layers of atoms with a periodic stacking sequence. In a tiling model, the layers are constructed by decorating the tiles of a quasiperiodic tiling – usually the Penrose rhomb tiling or one of its variants such as the Penrose kite and dart pattern (Gardner, 1977; Grünbaum & Shephard, 1987) or the first quasiperiodic pattern discovered by Penrose that contains regular pentagonal tiles (Penrose, 1974). These various Penrose tilings are all equivalent in the sense that there are ‘recombination’ rules which will convert one variety to another; moreover, each variety of Penrose tiling is equivalent to itself with a scaling factor \( \tau = (1 + \sqrt{5})/2 \) (Penrose, 1978; Grünbaum & Shephard, 1987; Lord, 1991). Note, however, that decorating the set of tiles of a quasiperiodic tiling does not in general give rise to a consistent decoration of the tiles of an equivalent tiling. (For example, consider a Penrose tiling consisting of white fat rhombs and black thin rhombs; in the associated kite and dart tiling, the kites are not all identically marked, nor are the darts.)

Gummelt’s discovery (Gummelt, 1995a,b) of a quasiperiodic covering of the plane using a single ‘tile’ has been exploited recently to produce a convincing model for the decagonal phase of Al–Ni–Co (Steinhardt et al., 1998). In Gummelt’s scheme, the tiles are regular decagons that overlap rather than meeting edge to edge. The matching rule requires that a black and white decoration, identical for every decagon, shall be consistent on the overlap regions (black on black, white on white). Steinhardt et al. elucidated the structure of decagonal Al–Ni–Co in terms of an atomic decoration of the Gummelt tile, which in 3D is a decagonal prism. This decorated decagonal prism is analogous to a unit cell in a periodic structure, except that the units overlap, sharing atoms at ‘coincidence sites’ in the overlap regions. The transformations relating pairs of contiguous clusters are reflections and glide reflections rather than translations. Steinhardt et al. named this basic unit a ‘quasi unit cell’.

The concept introduced by Steinhardt et al. for the elucidation of this particular phase raises a more general question. What characteristics of a quasiperiodic pattern in two or in three dimensions are necessary and sufficient for it to be describable in terms of a decagonal quasi unit cell, and how can these patterns be classified?

2. Gummelt coverings in 2D

Gummelt showed that every covering of the plane by Gummelt decagons can be converted to a unique kite and dart (KD) tiling by superimposing a ‘cartwheel’ on every decagon, as in Fig. 1.

A simpler relation between the kite and dart tilings and the Gummelt coverings is indicated in Fig. 2 (Lord et al., 2000). The patch consisting of a dart and two kites was called an ‘ace’ by Conway, who showed that every tile of a kite and dart tiling
belongs to an ace (Grünbaum & Shephard, 1987). The figure indicates how every Gummelt decagon can be converted to an ace and *vice versa*. The Gummelt coverings and the Penrose KD tilings are in this sense equivalent. The ‘cartwheel’ of Fig. 1 is obtained by applying two ‘decompositions’ (Grünbaum & Shephard, 1987; Lord, 1991) to the dart and kites of the ace in Fig. 2. Contiguous aces (*i.e.* aces abutting along an edge or sharing a kite) of a kite and dart tiling or, *equivalently*, pairs of overlapping decagons in a Gummelt covering, are related by one of the following transformations:

- **A**: \(4\pi/5\) (anticlockwise) about the point marked \(a, b\);
- **B**: \(2\pi/5\) about \(a, b\);
- **C**: \(2\pi/5\) about \(c\);
- **D**: \(\pi/5\) about \(d\);
- **E**: \(\pi/5\) about \(e\);
- **A\(^{-1}\)**, **B\(^{-1}\)**, **C\(^{-1}\)**, **D\(^{-1}\)**, **E\(^{-1}\)**.

In Fig. 3, we have illustrated the sets of decagons obtained by applying the cyclic groups of transformations generated by **B** (or **A** \(\hat{=}\) **B**\(^2\)), by **C** and by **D** (the figure generated by **E** would of course be just the mirror image of the **D** figure). Notice, incidentally, that the ring of ten decagons generated by **D** cannot actually occur in a covering of the plane satisfying Gummelt’s matching rules; the central decagonal hole cannot be covered without violating Gummelt’s rules; recall, in particular, the rule that every decagon edge must be contained in an overlap region (Gummelt, 1995\(^a\), \(^b\)). The configuration is the analogue, in the context of Gummelt coverings, of a ‘decapod’ defect in a Penrose tiling (Grünbaum & Shephard, 1987, p. 267).

Our aim is to consider the aperiodic patterns, in two and three dimensions, that can be constructed by decorating the decagons of a Gummelt covering, and to classify the possibilities. The symmetries of the portion of a pattern that lies in the overlap regions of the patches shown in Fig. 3 are the basis of our proposed classification scheme.

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### Figure 2
An ace (the dart and the two kites outlined by heavy lines) superimposed on a Gummelt decagon and the centres of rotations that relate pairs of contiguous aces or, equivalently, pairs of overlapping decagons.

### Figure 3
Patches generated by applying cyclic rotation groups to the Gummelt decagon. Overlap regions are outlined in heavier lines.
‘decorated’ identically by the portion of the pattern that falls within them, will be referred to as a pattern generated by the tiling or covering.

We use the word ‘pattern’ here in a very general sense. It may be a set of atoms of various kinds at various positions, or a function such as electron density, or simply an arrangement of arbitrary ‘shapes’. Periodic patterns in the plane, for example, are generated by tilings; the tiles are the fundamental regions or ‘asymmetric units’ for the appropriate ‘wallpaper group’ (Coxeter, 1961). Similarly, three-dimensional triply periodic patterns are generated by decorating the asymmetric units (as given in International Tables for Crystallography, 1987) of the space groups. The ‘matching rules’ that produce the tiling are the generators of the space group, which relate contiguous pairs of tiles. [See Lord (1997), in which continuous surfaces are derived as patterns generated by decorated asymmetric units.] The Penrose tilings have been employed extensively in the construction of models for decagonal quasicrystalline phases. An early example of an icosahedral quasicrystalline structure generated by an Ammann tiling (the 3D analogue of the Penrose rhomb tilings, consisting of two kinds of rhombic hexahedra) is the structure of icosahedral Al–Mn proposed by Yamamoto & Hiraga (1988).

A pattern generated by a Gummelt covering will be referred to as a G pattern.

Aperiodic patterns generated by decorated Penrose tilings or Gummelt coverings can be classified in a manner strikingly analogous to the space-group classification of periodic patterns. The classification we propose is based on the point symmetries of the decorations of the overlap regions of the clusters shown in Fig. 3. The portion of a pattern generated by a Gummelt covering (briefly: a ‘G pattern’) within a single decagon may have local reflection symmetries, as indicated by the lines of type X and Y in Fig. 4. (By a ‘local’ symmetry, we mean a symmetry that applies only to the portion of the pattern within the overlap regions indicated in Fig. 3. Reflection in a line Y, for example, relates a pair of decagons in the non-connected overlap region D of Fig. 3 but is not to be regarded as applying to the rest of the decagon.) We get just three kinds of plane G patterns, listed in Table 1. A quasi unit cell of a pattern of a given type is obtained by placing an object at a general position in the decagon and then applying to it the symmetry operations indicated in the corresponding row of the table.

4. The DKL tilings

Although a Gummelt covering is equivalent to a Penrose kite and dart tiling or a Penrose rhomb tiling, the equivalence does not in general carry over to the corresponding generated patterns. There are G patterns that do not have all the kites (or all the darts) of Fig. 1 decorated identically when the cart-wheel pattern is overlaid on all the decagons of the covering. There is, however, a tiling closely related to the KD tiling whose decorations give precisely all the G patterns.

In Fig. 5, a Gummelt decagon has been tiled by three kinds of tile: darts (D), kites (K) and ‘large kites’ (L). Using the set of transformations A, B, C, D, E and their inverses, a decoration of the decagon has the required consistency properties if and only if all the tiles of each of the three sets are identically decorated. In case the mirror symmetries Y are present (but not X), then the decorated tiles in the upper part of the figure are mirror images of corresponding tiles in the lower part. When the mirror symmetries X are present, each tile can be further subdivided, by mirror lines, into a pair of identical isosceles triangles. We shall refer to these triangular tiles as D, K and L triangles.

Table 1

<table>
<thead>
<tr>
<th></th>
<th>X</th>
<th>Y</th>
<th>D, E</th>
<th>B, C</th>
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<tr>
<td>p10</td>
<td>–</td>
<td>–</td>
<td>10</td>
<td>5</td>
</tr>
<tr>
<td>p5m</td>
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<td>p10m</td>
<td>m</td>
<td>m</td>
<td>10m</td>
<td>5m</td>
</tr>
</tbody>
</table>

Figure 4

The two kinds of local reflection symmetries of G patterns.

Figure 5

Darts, kites and ‘large kites’ superimposed on a Gummelt decagon. The ‘large kites’ are the white regions.
The ‘matching rules’ for building up DKL tilings can be expressed in terms of three kinds of vertex markings, as indicated in Fig. 6(a) in terms of DKL triangles.

A DKL tiling can be produced from a KD tiling by the rule: convert every ace that is related to contiguous aces by the transformations $D$ and $E^{-1}$ to a large kite. This suggests immediately an obvious way of producing a decomposition rule for the DKL tilings: convert the given DKL large kites to aces, apply the KD decomposition rule, then convert the resulting $\tau^{-1}$ scaled KD tiling to DKL by the above rule. This decomposition rule does not satisfy the criterion that all the original tiles of each kind are dissected in an identical way — some of the K tiles become L tiles while others become aces (D + 2K). Thus the $\tau$ inflation rule for the KD tilings does not give rise to a consistent $\tau$ inflation rule for the associated Gummelt coverings. However, applying two successive decompositions gives the consistent $\tau^{-2}$ decomposition of D, K and L triangles shown in Fig. 6(b). Correspondingly, we get a $\tau^2$ inflation rule for the Gummelt coverings.

The DKL tiles are analogous, for the decagonal patterns, to the fundamental regions (or ‘asymmetric units’) of wallpaper patterns (or, in 3D, of the space groups). Special positions relative to these tiles are analogues of Wyckoff positions. The tile vertices are of three kinds, as in Fig. 6(a).

From the $\tau^{-2}$ decomposition rule, we have

$$
\begin{align*}
D & \rightarrow D + K + L \\
K & \rightarrow 2D + 3K + L \\
L & \rightarrow 5D + 7K + 3L,
\end{align*}
$$

from which we deduce the ratios of the numbers of tiles of each kind (in a ‘sufficiently large’ patch). We have

$$
\begin{align*}
n_D & : n_K : n_L = \tau : 5^{1/2} : 1.
\end{align*}
$$

Whenever a decagonal phase is realised as a G pattern and atomic positions in the ‘quasi unit cell’ are known or postulated, this formula provides a very simple way of computing atomic percentages.

5. Classification of three-dimensional G patterns

In the three-dimensional G patterns, the basic unit is a decagonal prism, the DKL tiles are prisms and each of the three types of pattern listed in Table 1 leads to several distinct classes of three-dimensional patterns, corresponding to the fact that the rotations $A$, $B$, $C$, $D$ and $E$ may be realised as screw transformations and the reflections $X$ and $Y$ may be realised as c glides or as diad rotations about ‘horizontal’ axes (i.e. perpendicular to the prism axis – the periodic axis). There may also be ‘horizontal’ mirror planes. We consider only cases in which the nature of the transformation $X$ is the same for all three tile types (the generalization – in which the very concept of ‘pattern’ becomes obscure – seems unlikely to lead to anything of interest).

The resulting classification of pentagonal and decagonal G patterns is summarized in Table 2. We shall define two G patterns as belonging to the same class, or as being of the same ‘type’, if their local symmetries $A$ and $B$ are the same. The analogy with the space-group classification of tetragonal and hexagonal periodic patterns is quite striking, though there are

<table>
<thead>
<tr>
<th>$X$</th>
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</tr>
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</table>
some subtle variations. The names we have given to the classes of \( G \) patterns (column 1) should not be confused with the usual symmetry classification of decagonal quasicrystals. The usual ‘symmetry’ classification of decagonal quasicrystalline phases operates in terms of five-dimensional space groups. The classification of \( G \) patterns given in Table 2 makes use only of considerations in three-dimensional space, and is able to indicate finer features – note for example that fivefold screw axes \( 5\alpha \) and tenfold screw axes \( 10\beta \) can occur in the same pattern, with \( p \) and \( q \) independent. On the other hand, decagonal quasicrystals are not necessarily ‘\( G \) patterns’ – probably most are not. The table identifies 165 possible types of three-dimensional \( G \) patterns. It is unlikely that more than a very few of them are realised in nature as quasicrystalline structures. Indeed, the centrosymmetric classes \( P10_5/mmc \) and \( P10/mmm \) occur predominantly as ‘symmetries’ of decagonal phases. [A non-centrosymmetric phase of Al–Mn–Pd with five-dimensional space group \( P10_5mc \) has been described in detail by Weber & Yamamoto (1998). It is not a \( G \) pattern.]

In the following section, we demonstrate how \( P10_5/mmc \) arises very naturally for \( G \) patterns or modified \( G \) patterns built from towers of icosahedra, ‘double icosahedra’ or pentagonal antiprisms.

6. Structure of some decagonal phases

In this section, we examine four decagonal quasicrystal structures in order to demonstrate that they can be regarded as \( G \) patterns or as systematic modifications of \( G \) patterns.

Fig. 7 shows a Gummelt decagon superimposed on a patch of three pentagonal tiles – indicating an obvious sense in which the Penrose pentagon tilings and the Gummelt coverings are equivalent. The three kinds of vertices introduced in Fig. 6 are here labelled \( a \), \( b \) and \( c \), and the vertices of the pentagons, \( d \), are special positions for the \( G \) patterns. In three-dimensional \( G \) patterns, positions of types \( a \), \( b \) and \( c \) are related to those of types \( a \), \( b \) and \( c \) by the ‘\( Y \) transformations’.

A very simple prescription converts a Penrose rhomb tiling satisfying the usual matching rules to an equivalent tiling consisting of three tile types: a ‘hex’ built from one fat rhomb and two thin rhombs, a ‘boat’, built from three fat and one thin, and a ‘star’ built from five fat rhombs (Tang & Jaric, 1990; Henley, 1991; Li & Kuo, 1991; Lord, 1991; Li, 1995; Cockayne & Widom, 1998b). The positions \( a \) of a Gummelt covering are the vertices of an HBS (hex, boat, star) tiling.

Rather surprisingly, the structure of decagonal Al–Mn, presented by Li (1995) in terms of a decorated HBS tiling, turns out to be a \( G \) pattern at a smaller scale. The basic structure of this phase consists of rods, or ‘towers’, of Al icosahedra (with centred Mn atoms) and pentagonal bipyramids (Fig. 8). If regularity of these polyhedra is assumed, the five layers labelled 0, 1, 2, 3, 4, 5 have \( z \) coordinates given approximately by 0.0, 0.11, 0.21, 0.29, 0.39 and 0.50. Three of the towers appear in projection as the small decagons in Fig. 7. The structure can be completely described in terms of ‘Wyckoff’ positions:

\[
\begin{align*}
\text{Al} : & \quad a(0, 0.39), b(0.29), c(0.11), d(0, 1/2) \\
\text{Mn} : & \quad a(0.29), c(1/2).
\end{align*}
\]

Fig. 9 indicates the decoration of the fundamental units (DKL triangular prisms). (Observe, incidentally, that the positions of the atoms on the sides of these prisms enforce the DKL matching rule!) Denoting...
the period by \( c \) and the edge of the decagonal quasi unit cell (Fig. 7) by \( a \), then \( c/a = 3.53 \). The HBS tiles employed by Li are quite large in relation to the decagonal cell. The edge length of the HBS tiles is actually \( r a \). Therefore, although Li’s model can, obviously, be described as a decorated Penrose rhomb tiling by partitioning the HBS tiles, the description in terms of a decorated DKL tiling – and the equivalent description as a ‘\( G \) pattern’ – is very much simpler. The ratio \( \text{Al}:\text{Mn} \) is readily computed from Fig. 9. The content of the prisms is \( D: (1/2)\text{Al} + (1/5)\text{Mn}, K: (11/20)\text{Al} + (1/5)\text{Mn}, L: (3/2)\text{Al} + (1/4)\text{Mn} \). Together with the ratio \( D:K:L = \tau:5^{1/2}:1 \), this gives

\[
\text{Al}: \text{Mn} = 10\tau + 11 \times 5^{1/2} + 30 : 4\tau + 4 \times 5^{1/2} + 5 = 3.45 : 1.
\]

All the sites in the quasi unit cell of the Al–Ni–Co model of Steinhardt et al. (1998) are given by the \( a, b \) and \( c \) positions of a \( G \) pattern based on a smaller decagon than the quasi unit cell (smaller by a factor \( \tau^{-1} \)). This is illustrated in Fig. 10. The positions \( a \) of the \( \tau^{-1} \) decagons (which coincide with the positions \( c \) of the quasi unit cell) are occupied by Al, \( b \) by Al and Co, \( c \) by Al, Ni and Co. A few of the \( b \)-type positions are vacant. These vacancies are necessary for the exact \( G \) pattern of the Al–Ni–Co model – an anomaly related to the anomalous \( \tau \) inflation for the Gummelt coverings that we have already alluded to in our discussion of DKL tilings. The existence of this \( \tau^{-1} \) scaled \( G \) pattern implies that the structure of decagonal Al–Ni–Co is basically an arrangement of towers of pentagonal antiprisms (appearing in projection as the small decagons of \( b \) and \( c \) sites in Fig. 7). Assuming equal edge lengths, it is not difficult to deduce that \( c/a = \tau^{-1}5^{-1/4}\times \cos(18°) = 0.384 \) (where \( a \) is the edge length of the actual decagonal quasi unit cell, not the \( \tau^{-1} \) scaled cell).

The decoration of the quasi unit cell for Al–Ni–Co almost corresponds to the decoration of a DKL tiling shown in Fig. 11. Positions \( a(1/2), b(0) \) and \( c(0) \) are occupied, respectively, by cobalt, nickel and aluminium atoms. Additional Al and Co sites are at special positions of the \( \tau^{-1} \) pattern, as described above. However, the cobalt atom and the two aluminium atoms that form a small triangle inside the large kites are absent in the large kite at the centre of the quasi unit cell. These sites are occupied only when this large kite is overlapped by occupied sites. (This subtlety is a deviation from the strict concept of \( G \) pattern as we have defined it.) The calculation of the ratio \( \text{Al}:\text{Ni}:\text{Co} \) is consequently not so straightforward as the computation of \( \text{Al}:\text{Mn} \) from Li’s model. This problem is dealt with in Appendix A.

The decagonal Al–Cu–Co phase elucidated by Cockayne & Widom (1998a) is a decoration of an HBS tiling. Wittmann (1999) has already given a detailed discussion of the deviations of this structure from an exact \( G \) pattern of the kind proposed by Steinhardt et al. for Al–Ni–Co. We approach this question from a somewhat different viewpoint, emphasizing the \( \tau^{-1} \) \( G \) pattern of sites and the towers of pentagonal antiprisms. Consider a decagonal quasi unit cell with \( a, b \) and \( c \) positions occupied as follows:

\[
\text{Co} : a(3/4); \text{Al}/\text{Cu} : b(1/4); \text{Al} : c(0, 1/2).
\]

Additional sites introduced by the \( \tau^{-1} \) \( G \) pattern are

\[
\text{Al} : a(0, 1/2), b(3/4); \text{Al}/\text{Cu} : b(1/4).
\]

In the model of Cockayne & Widom, the distribution of the copper atoms at the Al/Cu positions is not random but is mainly determined by the requirement that the edges of the

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**Figure 9**
Decoration of the DKL prisms that reproduces Li’s decagonal Al–Mn. Open circles denote Al, black circles Mn. The numbers refer to the layers perpendicular to the periodic axis.

**Figure 10**
The \( a, b \) and \( c \) positions of a Gummelt covering by decagons of edge length \( \tau^{-1} \) are all the sites (including vacant sites) for the Al–Ni–Co structure with quasi unit cell of unit edge length. White and black circles correspond, respectively, to \( z = 0 \) and \( z = 1/2 \). Atoms marked with arrows are absent in large kites that belong to only one decagon.

**Figure 11**
Decoration of the DKL prisms that reproduces the Al–Ni–Co model of Steinhardt et al.
HBS tiling (whose vertices are the $c$ positions of the $\tau^{-1}$ pattern) each contain one Cu and one Co; this enforces the HBS matching rule.

The deviations of the structure from an exact $G$ pattern are due to the particular distribution of Cu and the absence of vacancies at the Al/Cu sites, and the placement of Al atoms at $c(0, 1/2)$ (i.e. at the centres of the pentagonal antiprisms that surround them) rather than at $c(1/4)$.

Cockayne & Widom (1998b) have also made a detailed study of orthorhombic approximants to the decagonal Al–Co quasicrystalline phase. These approximants are decorations of periodic HB tilings, from which the corresponding decoration of the quasiperiodic HBS tiles are readily deduced.

The basic structure of the Al–Co model can be visualized by interpreting the small decagons in Fig. 7 as projections of towers of `double icosahedra'. A view along the `Y axis' is illustrated in Fig. 12. Assuming equal bond lengths gives

$$c/a = (4/5)\cos(18^\circ)\tau^2(\tau + 2)^{1/2} = 0.857.$$  

An exact $G$ structure is

In the actual Al–Co phase as deduced from its approximants, additional ‘interstitial’ Al atoms occur at $c(1/2)$ and $c'(0)$. The short diagonal of the D tiles is in fact too short to accommodate atoms at both its vertices in the same layer. The interstitials thus cause a shift of the positions of neighbouring atoms and consequent deviations from an exact $G$ pattern. In particular, the fivefold symmetry of the towers is disrupted. The position shifts are indicated in Fig. 13.

7. Conclusions

The recognition of the existence of quasiperiodic tilings (the DKL tilings) in 2D and 3D whose decorations produce the $G$ patterns (i.e. the structures obtained by decorating the decagons of a Gummelt covering) leads to a simple way of labelling the important sites in a certain class of quasicrystals. The identification of four known $P10_5/mmc$ quasicrystals as Gummelt structures or systematically disrupted Gummelt structures suggests that this approach may be worth further investigation.

APPENDIX A

The structure of decagonal Al–Ni–Co proposed by Steinhardt et al. does not strictly fit our definition of a ‘$G$ pattern’. In a Gummelt covering, the central region of a decagon (white L tile in Fig. 5) may belong to more than one decagon or to only one. In the Al–Ni–Co model, three atoms (two Al and one Co) are replaced by vacancies in those L tiles that are not shared. However, a single quasi unit cell with three vacancies in the middle is sufficient to define the structure; the vacant sites become occupied whenever they coincide with atoms from another quasi unit cell. Thus, in order to obtain the ratios Al:Ni:Co predicted by the model we need to know what proportion of the decagons in a Gummelt covering have central L tiles that are not shared by any other decagon. The solution of this problem presented here calls into play some interesting geometrical properties that are not without interest in their own right.

Observe that, since an HBS tiling is very simply obtained by omitting certain edges in a Penrose rhomb tiling (specifically, those with double arrow heads, in the conventional edge labelling that enforces the matching rule), a $\tau^{-1}$ deflation rule for HBS tilings is readily obtained from that of the rhomb tilings. Deflating twice gives the $\tau^{-2}$ rule shown in Fig. 14,

$$H \rightarrow 3H + B + S$$

$$B \rightarrow 4H + 3B + S$$

$$S \rightarrow 5H + 5B + S,$$

from which it follows that

$$H \rightarrow 3H + B + S$$

$$B \rightarrow 4H + 3B + S$$

$$S \rightarrow 5H + 5B + S,$$
This gives (in the limit, for a sufficiently large patch of HBS tiling)

$$n_H : n_B : n_S = \tau^{5/2} : 5^{1/2} : 1.$$
Referring to Fig. 15, we see that each H contributes $2L_0 + 9L_1$, each B contributes $3L_0 + 15L_1$, and each S contributes $5L_0 + 20L_1$. Since $n_H:n_0:n_S = \tau:5^{1/2}:1$, we have

$$n_0 : n_1 = 2\tau 5^{1/2} + 3 \times 5^{1/2} + 5 : 9\tau 5^{1/2} + 15 \times 5^{1/2} + 20 = 8\tau + 39 : 23,$$

Normalizing so that $n_0 + n_1 = 1$,

$$nD : nK : n_0 : n_1 = \tau : 5^{1/2} : 10\tau - 16 : 17 - 10\tau.$$

The numbers of atoms of each kind, contributed by the D, K, L_0 and L_1 tiles (see Fig. 11), are:

<table>
<thead>
<tr>
<th></th>
<th>Al</th>
<th>Ni</th>
<th>Co</th>
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<tr>
<td>D</td>
<td>9/5</td>
<td>1/5</td>
<td>0</td>
</tr>
<tr>
<td>K</td>
<td>7/5</td>
<td>2/5</td>
<td>1/5</td>
</tr>
<tr>
<td>L_0</td>
<td>13/5</td>
<td>7/5</td>
<td>0</td>
</tr>
<tr>
<td>L_1</td>
<td>23/5</td>
<td>7/5</td>
<td>1</td>
</tr>
</tbody>
</table>

Finally, one obtains the ratio

$$\text{Al} : \text{Ni} : \text{Co} = 9\tau + 7 \times 5^{1/2} + 13(10\tau - 16) + 23(17 - 10\tau) : \tau + 5^{1/2} + 7 : 5^{1/2} + 5(17 - 10\tau)$$

$$= 176 - 77\tau : 5\tau + 5 : 84 - 48\tau,$$

which gives $72.6 : 18.5 : 8.9$ [in close agreement with the ratios given in Steinhardt et al. (1998)].

An alternative approach to the same problem is not without interest; we shall indicate it briefly. Instead of the pentagon vertices given by the centres of the decagonal cells, consider the pentagon vertices given by all the d sites (see Fig. 7), which lie inside the L tiles. In Fig. 16, those inside L_0 tiles are marked by white circles, those inside L_1 tiles by grey or black circles. (The white and black circles are at centres of the decagonal quasi unit cells.) This information can be verified by constructing the actual Gummelt covering that Fig. 16 encodes. Looking now at the associated HBS tiling, we see that each H contributes 4 L_1, each B contributes L_0 + 6 L_1 and each S contributes 5L_0 + 5 L_1. Hence, $n_0:n_1 = 5^{1/2} : 5\tau 5^{1/2} + 6\tau + 5 = 2\tau + 4 : 16\tau + 7$. Again, normalizing to $n_0 + n_1 = 1$ gives $n_0:n_1 = 10\tau - 16 : 17 - 10\tau$.

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**References**


